

On the regularity of null-controls of the linear 1-d heat equation

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Abstract

The $1 - d$ heat equation in a bounded interval is null-controllable from the boundary. More precisely, for each initial data $y^0 \in L^2(0, 1)$ it corresponds a unique boundary control of minimal $L^2(0, T)$ -norm which drives the state of the $1 - d$ linear heat equation to zero in time $T > 0$. This control is given as the normal derivative of a solution of the homogeneous adjoint equation whose initial data $\hat{\varphi}^0$ minimizes a suitable quadratic cost functional. In this Note we analyze the relation between the regularity of the initial datum to be controlled y^0 and that of $\hat{\varphi}^0$. We show that, if y^0 has only one Fourier mode, the corresponding $\hat{\varphi}^0$ does not belong to any Sobolev space of negative exponent. This explains the severe ill-posedness of the problem and the lack of efficiency of most of the existing numerical algorithms for the numerical approximation of the controls, and in particular the slow convergence rate of Tychonoff regularization procedures.

Sur la régularité des contrôles pour l'équation de la chaleur linéaire

Résumé

L'équation de la chaleur $1 - d$ dans un intervalle borné est contrôlable depuis le bord. Plus précisément, pour tout $y^0 \in L^2(0, 1)$ il existe un unique contrôle frontière de norme $L^2(0, T)$ minimale qui conduit la solution de l'équation linéaire de la chaleur à zéro à l'instant $T > 0$. Ce contrôle est donné par la dérivée normale au bord d'une solution de l'équation adjointe avec une donnée initiale $\hat{\varphi}^0$ qui minimise une fonction coût quadratique. Dans cet article nous étudions la relation entre la régularité de y^0 et celle de $\hat{\varphi}^0$. Nous montrons que, si y^0 a une seule fréquence de Fourier, la fonction $\hat{\varphi}^0$ correspondante n'appartient à aucun espace de Sobolev d'exposant négatif. Ce fait explique le caractère mal posé du problème et l'inefficacité de la plupart des algorithmes numériques existants et en particulier la convergence lente des méthodes de régularisation de type Tychonoff.

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Version française abrégée

L'équation de la chaleur 1 - d dans un intervalle borné (1) est contrôlable depuis le bord. Plus précisément, pour tout $y^0 \in L^2(0, 1)$, il existe un unique contrôle frontière v de norme $L^2(0, T)$ minimale qui conduit la solution y de l'équation linéaire de la chaleur à zéro à l'instant $T > 0$: $y(T) \equiv 0$.

Ce contrôle est donné par la dérivée normale au bord d'une solution de l'équation adjointe (2) avec une donnée initiale $\tilde{\varphi}^0$ qui minimise la fonction coût quadratique \mathcal{J} donnée par (3). Le minimum de cette fonction est atteint sur un espace de Hilbert ad-hoc \mathcal{H} défini à partir de l'inégalité d'observabilité (4) du système adjoint.

Cet espace est difficile à caractériser en termes classiques. Il est connu qu'il est plus large que tout espace de Sobolev d'indice négatif à cause du fort phénomène régularisant de l'équation de la chaleur. Ce fait est lié au caractère mal posé du problème au sens que, en particulier, les méthodes numériques usuelles ont souvent des vitesses de convergence très lentes ([2,4,8]). C'est, par exemple, le cas des méthodes d'approximation basées sur la régularisation de Tychonoff.

Dans cette Note nous donnons une réponse précise à cette question en démontrant que les minima de la fonctionnelle \mathcal{J} sont souvent aussi irréguliers que l'espace \mathcal{H} le permet. Nous faisons cela à travers une analyse soignée de la relation entre la régularité de y^0 et celle de $\tilde{\varphi}^0$. Nous montrons, en particulier, que, si y^0 a une seule fréquence de Fourier, la fonction $\tilde{\varphi}^0$ correspondante n'appartient à aucun espace de Sobolev d'exposant négatif car ses coefficients de Fourier croissent de manière exponentielle à hautes fréquences. Notre preuve est basée sur une analyse des minimiseurs à l'aide des estimations fines sur la suite de fonctions biorthogonales des exponentielles intervenant dans le développement en série de Fourier des solutions.

Ceci donne une explication définitive au caractère mal posé du problème et à l'inefficacité de la plupart des algorithmes numériques existants et en particulier la lente convergence des méthodes de régularisation de type Tychonoff qui, pour établir une vitesse de convergence, ont besoin d'un minimum de régularité Sobolev des minima.

Notre analyse s'applique, avec des résultats analogues, lorsque le contrôle agit sur un point à l'intérieur du domaine ou lorsqu'il est tout simplement l'intensité d'un profil de contrôle donné distribué en espace.

1. Introduction

Given $T > 0$, $y^0 \in L^2(0, 1)$ and $v \in L^2(0, T)$, we consider the non-homogeneous 1-d heat equation

$$\begin{cases} y_t(t, x) - y_{xx}(t, x) = 0 & x \in (0, 1), t \in (0, T) \\ y(t, 0) = 0, \quad y(t, 1) = v(t) & t \in (0, T) \\ y(0, x) = y^0(x) & x \in (0, 1). \end{cases} \quad (1)$$

System (1) is well known to be *null-controllable (or controllable to zero) in time T* for all $T > 0$. More precisely, for each $y^0 \in L^2(0, 1)$ there exists a *control* $v \in L^2(0, T)$ such that the corresponding solution of (1) verifies $y(T, \cdot) = 0$. The control of minimal $L^2(0, T)$ -norm is unique and can be characterized through the adjoint system as described below.

Given $\varphi^0 \in L^2(0, 1)$ we consider the adjoint heat equation

$$\begin{cases} \varphi_t(t, x) + \varphi_{xx}(t, x) = 0 & x \in (0, 1), t \in (0, T) \\ \varphi(t, 0) = \varphi(t, 1) = 0 & t \in (0, T) \\ \varphi(T, x) = \varphi^0(x) & x \in (0, 1). \end{cases} \quad (2)$$

In view of the regularizing properties of the heat equation, the map $\varphi^0 \rightarrow \left(\int_0^T (\varphi_x)^2(t, 1) dt \right)^{\frac{1}{2}}$ is a norm in $L^2(0, 1)$. Let \mathcal{H} be the completion of $L^2(0, 1)$ with respect to this norm. Now, we introduce the functional $\mathcal{J} : \mathcal{H} \rightarrow \mathbb{R}$ given by

$$\mathcal{J}(\varphi^0) = \frac{1}{2} \int_0^T |\varphi_x|^2(t, 1) dt - \int_0^1 y^0(x) \varphi(0, x) dx, \quad (3)$$

φ being the solution of (2) with initial data φ^0 .

For all $y^0 \in L^2(0, 1)$, \mathcal{J} is coercive in \mathcal{H} . This is a deep result that requires of the following *observability inequality*: For all $T > 0$ there exists $C > 0$ such that

$$\|\varphi(0)\|_{L^2(0,1)}^2 \leq C \int_0^T |\varphi_x(1, t)|^2 dt, \quad \forall \varphi^0 \in L^2(0, 1). \quad (4)$$

This kind of result is by now well known also in the multi-dimensional case. We refer to [9] for an updated account on the developments in this context.

The functional \mathcal{J} being coercive in \mathcal{H} , in addition to being continuous and strictly convex, it has a unique minimizer $\hat{\varphi}^0 \in \mathcal{H}$. The normal derivative of the solution of the adjoint system associated to this minimizer yields the control of minimal $L^2(0, T)$ -norm. Summarizing, the following holds (see, for instance, [6]).

Proposition 1 *Let $T > 0$ be given. For each $y^0 \in L^2(0, 1)$ there exists a unique control $v \in L^2(0, T)$ of minimal $L^2(0, T)$ -norm for equation (1) and $v = \hat{\varphi}_x(\cdot, 1)$, where $\hat{\varphi}$ is the solution of the adjoint problem (2) with initial data $\hat{\varphi}^0 \in \mathcal{H}$, the minimizer of (3). Moreover, the map $\mathcal{G} : L^2(0, 1) \rightarrow \mathcal{H}$, $\mathcal{G}(y^0) = \hat{\varphi}^0$, is linear and continuous.*

Despite of this clear functional description and characterization of this optimal null control, as reported in [2,4,8], the problem of computing $\hat{\varphi}^0 \in \mathcal{H}$ by minimizing (3) is severely ill-posed. This comes from the fact that the space \mathcal{H} is very large. In fact, due to the regularizing effect of the heat equation, one can see that any distribution in a negative order Sobolev space $H^{-s}(0, 1)$, with support away from $x = 1$, belongs to \mathcal{H} , whatever $s > 0$ is. Therefore, for a given $y^0 \in L^2(0, 1)$, the minimizer of \mathcal{J} may have very low regularity and it may be difficult to capture it numerically with accuracy and robustness.

The main aim of this article is to investigate the regularity of $\mathcal{G}(y^0) = \hat{\varphi}^0 \in \mathcal{H}$ and to show that, indeed, for regular initial data y^0 , as it is for instance the case when y^0 is constituted by one single Fourier mode, the corresponding minimizer $\mathcal{G}(y^0)$ of \mathcal{J} does not belong to any Sobolev space of negative exponent.

This clearly expresses the severe ill-posedness of the problem and also shows why classical methods, based on Tychonoff regularization (see [5]), fail to be efficient. Indeed, the existing results about the convergence rates for Tychonoff regularization require a minimal regularity of the minimizer $\mathcal{G}(y^0) = \hat{\varphi}^0$ of \mathcal{J} in some Sobolev space, something that, as described above, is not fulfilled.

The spectral approach we have developed in this paper allows also handling other types of controls. Indeed, instead of looking for a boundary control v as in (1), we may be consider an interior ‘‘lumped

control” of the form $v(t)g(x)$, where $g \in L^2(0,1)$ is a given profile control function, or a pointwise control like $v(t)\delta_{x_0}$, where δ_{x_0} is the Dirac mass concentrated at some point $x_0 \in (0,1)$ (see, for instance, [1,4]). The analysis of the family of biorthogonal functions remains the same. The only difference when identifying the datum of the adjoint system at the minimizer is that its Fourier coefficients are affected by the weights that these control mechanisms introduce in the observability inequalities. However, our results, in the sense that the Fourier coefficients of the optimal adjoint state grow exponentially, remain true.

2. Analysis of the operator \mathcal{G}

Let us introduce some notation. The eigenvalues of the 1-d Dirichlet Laplace operator are $\lambda_n = n^2\pi^2$ and the corresponding eigenfunctions $\Phi^n = \sin(n\pi x)$, for every $n \in \mathbb{N}^*$. $\Lambda = (e^{-\lambda_n t})_{n \geq 1}$ denotes the family of the corresponding real exponential functions and, for any $T > 0$, $E(\Lambda, T)$ is the closure of the algebraic span of Λ in $L^2(0, T)$. It is known that, for any $T > 0$, there exists a unique biorthogonal family $(\theta_T^m)_{m \geq 1}$ to Λ in $L^2(0, T)$ such that $(\theta_T^m)_{m \geq 1} \subset E(\Lambda, T)$. Moreover, this biorthogonal sequence has minimal $L^2(0, T)$ -norm.

Let v be the control of minimal $L^2(0, T)$ -norm corresponding to y^0 and $\widehat{\varphi}^0 = \mathcal{G}(y^0)$ given by Proposition 1. As we have said before, our aim is to study the regularity of $\widehat{\varphi}^0$ and we do this by expanding it in Fourier series

$$\mathcal{G}(y^0) = \widehat{\varphi}^0(x) = \sum_{m \geq 1} b_m \sin(\pi m x) \quad (5)$$

and analyzing the behavior of the Fourier coefficients b_m that can be characterized as follows.

Theorem 2.1 *Let $y^0 = \sum_{n \geq 1} a_n \sin(\pi n x) \in L^2(0, 1)$. The Fourier coefficients b_m from (5) are given by*

$$b_m = \frac{(-1)^m}{m\pi} \sum_{n=1}^{\infty} (-1)^n \frac{a_n}{2n\pi} e^{-n^2\pi^2 T} d_n^m \quad \forall m \geq 1 \quad (6)$$

where $d_n^m = \langle \theta_T^n, \theta_T^m \rangle_{L^2(0, T)}$ for all $m, n \geq 1$.

Proof: From the minimal norm property of θ_T^n in $L^2(0, T)$, we deduce that the control of minimal $L^2(0, T)$ -norm corresponding to y^0 is given by the formula

$$v(t) = \sum_{n=1}^{\infty} (-1)^n \frac{a_n}{2n\pi} e^{-n^2\pi^2 T} \theta_T^n(T-t). \quad (7)$$

If $\widehat{\varphi}$ is the solution of (1) with initial data $\widehat{\varphi}^0$ given by (5), $v = \widehat{\varphi}_x(\cdot, 1)$ and (6) follows immediately. ■

In order to evaluate the Fourier coefficients b_m of $\mathcal{G}(y^0)$ we need to estimate the quantities d_k^n . This is done in the following key result whose proof requires of significant technical developments that can be found in [6].

Theorem 2.2 *There exist a positive integer n_0 and a positive constant $C > 0$, independent of k and m , but depending of T , such that*

$$|d_k^m| \leq C e^{\pi(k+m)} \quad \forall k, m \geq 1 \quad (8)$$

$$|d_k^m| \geq \frac{k m}{32(k^2 + m^2)} e^{\pi(k+m)} \quad \forall m, k \geq n_0. \quad (9)$$

Sketch of the proof: We use a strategy similar to the one used in [3], where the norms $\|\theta_T^n\|_{L^2(0,T)} = \sqrt{d_n^n}$ are evaluated. It consists in “truncating” the infinite matrix $G = (d_k^m)_{k,m \geq 1}$ by considering the minimal norm biorthogonal $(\theta_\infty^m)_{1 \leq m \leq N}$ to the finite family $\Lambda_N = (e^{-\lambda_n t})_{1 \leq n \leq N}$ in $L^2(0, \infty)$. In this case we have an explicit formula for the quantities $d_k^m(N, \infty) = \langle \theta_\infty^m, \theta_\infty^k \rangle_{L^2(0, \infty)}$. Then, a perturbation argument allows to estimate $d_k^m(N, T) = \langle \theta_T^m, \theta_T^k \rangle_{L^2(0, T)}$, where $(\theta_T^m)_{1 \leq m \leq N}$ is the minimal norm biorthogonal to Λ_N in $L^2(0, T)$. Finally, by letting N to tend to infinity, we obtain the desired estimates for d_k^m . ■

3. Control theoretical consequences

In this section we deduce some consequences of Theorems 2.1 and 2.2 for the controllability of (1). Let us consider the case in which the initial data y^0 of (1) to be controlled is the n -th eigenfunction Φ^n of the Laplace operator. The following result is a direct consequence of (6) and estimates (8)-(9).

Corollary 3.1 *Let $n \geq n_0$ and $(b_m^{0,n})_{m \geq 1}$ be the Fourier coefficients of $\mathcal{G}(\Phi^n)$ from (5). Then*

$$|b_m^{0,n}| \geq \frac{1}{64\pi^2(n^2 + m^2)} e^{-\pi^2 T n^2 + \pi(m+n)} \quad \forall m \geq n_0. \quad (10)$$

The constants C and n_0 are those given by Theorem 2.2, are independent of n and m but depend of T .

Remark 1 Estimate (10) shows that, for any $n \geq n_0$, the initial data $\mathcal{G}(\Phi^n)$ of (2), which gives the control of minimal $L^2(0, T)$ -norm for the eigenfunction Φ^n , has a very low regularity. Indeed, for any $m \geq 2\pi T n^2$, we obtain that $|b_m^{0,n}| \geq e^{\frac{\pi}{2}m}$ and therefore $\mathcal{G}(\Phi^n)$ does not belong to any Sobolev space of negative order. □

The following result concerns the regularity of the control of minimal $L^2(0, T)$ -norm for arbitrary initial data in $L^2(0, 1)$.

Corollary 3.2 *For any $y^0 \in L^2(0, 1)$ and $\varepsilon > 0$ there exists $y_\varepsilon^0 \in L^2(0, 1)$ and a positive constant C , depending only of T and ε , such that $\|y^0 - y_\varepsilon^0\|_{L^2(0,1)} < \varepsilon$ and*

$$|b_{\varepsilon, m_k}^0| \geq \frac{C}{m_k^2} e^{\pi m_k} \quad \forall k \geq 1, \quad (11)$$

$(b_{\varepsilon, m}^0)_{m \geq 1}$ being the Fourier coefficients of $\mathcal{G}(y_\varepsilon^0)$ and $(m_k)_{k \geq 1}$ an increasing sequence of positive integers.

Proof: Let $\mathcal{G}(y^0) = \sum_{m=1}^{\infty} b_m^0 \sin(m\pi x)$ and n_0 be the positive integer from Theorem 2.2. We define $y^1 = y^0 + \varepsilon \sin(n_0 \pi x)$ and remark that $\mathcal{G}(y^1) = \mathcal{G}(y^0) + \varepsilon \mathcal{G}(\Phi^{n_0})$, where $\mathcal{G}(\Phi^{n_0})$ was studied in Corollary 3.1. If $(b_m^1)_{m \geq 1}$ are the Fourier coefficients of $\mathcal{G}(y^1)$, we deduce that

$$b_m^1 = b_m^0 + \varepsilon b_m^{0, n_0} = b_m^0 + \varepsilon \frac{(-1)^{n_0+m}}{2 n_0 m \pi^2} e^{-n_0^2 \pi^2 T} d_m^{n_0} \quad m \geq 1. \quad (12)$$

Consequently, $\max\{|b_m^0|, |b_m^1|\} \geq \frac{\varepsilon |d_m^{n_0}|}{4 n_0 m \pi^2} e^{-n_0^2 \pi^2 T}$ and, by taking into account (9) from Theorem 2.2, we deduce that at least one of the sequences $(b_m^0)_{m \geq 1}$ or $(b_m^1)_{m \geq 1}$ has a subsequence which verifies (11). If this is $(b_m^0)_{m \geq 1}$, we choose $y_\varepsilon^0 = y^0$. Otherwise, we take $y_\varepsilon^0 = y^1$ and the proof ends. ■

Remark 2 Corollary 3.2 shows that the set of initial data y^0 whose controls of minimal $L^2(0, T)$ -norm are given by minimizers $\mathcal{G}(y^0)$ of \mathcal{J} which do not belong to any Sobolev space of negative order is dense in $L^2(0, 1)$. \square

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